

UV and IR Spectroscopic Studies of Cold Alkali Metal Ion–Crown Ether Complexes in the Gas Phase

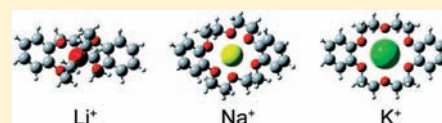
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 Supporting Information

ABSTRACT: We report UV photodissociation (UVPD) and IR-UV double-resonance spectra of dibenzo-18-crown-6 (DB18C6) complexes with alkali metal ions (Li^+ , Na^+ , K^+ , Rb^+ , and Cs^+) in a cold, 22-pole ion trap. All the complexes show a number of vibronically resolved UV bands in the $36\,000\text{--}38\,000\text{ cm}^{-1}$ region. The Li^+ and Na^+ complexes each exhibit two stable conformations in the cold ion trap (as verified by IR-UV double resonance), whereas the K^+ , Rb^+ , and Cs^+ complexes exist in a single conformation. We analyze the structure of the conformers with the aid of density functional theory (DFT) calculations. In the Li^+ and Na^+ complexes, DB18C6 distorts the ether ring to fit the cavity size to the small diameter of Li^+ and Na^+ . In the complexes with K^+ , Rb^+ , and Cs^+ , DB18C6 adopts a boat-type (C_{2v}) open conformation. The K^+ ion is captured in the cavity of the open conformer thanks to the optimum matching between the cavity size and the ion diameter. The Rb^+ and Cs^+ ions sit on top of the ether ring because they are too large to enter the cavity of the open conformer. According to time-dependent DFT calculations, complexes that are highly distorted to hold metal ions open the ether ring upon $S_1\text{--}S_0$ excitation, and this is confirmed by extensive low-frequency progressions in the UVPD spectra.



1. INTRODUCTION

Crown ethers have played important roles in host–guest chemistry for the selective capture of guest species, especially metal cations. In 1967, Pedersen discovered the first crown ether, dibenzo-18-crown-6 (DB18C6, Scheme 1), and investigated its complexation with metal salts^{1,2} using UV spectroscopy as one of three criteria to verify complex formation. Ultraviolet spectra of DB18C6 were measured in methanol at room temperature with various kinds of salts. It was found that cation–DB18C6 complexes gave absorption maxima at 273 and 279 nm, whereas the DB18C6 solution with no salt showed only one maximum at 274 nm.¹ Similar spectral evidence of complex formation was found for all the alkali metal ions with DB18C6.¹ Unfortunately, the broadening of the UV spectra from solvent and temperature effects prevented further information about the structure of the complexes from being derived. In the same paper, Pedersen used solubility measurements to demonstrate that among the alkali metal cations K^+ was the most strongly bonded to DB18C6.¹ Shizuka et al. reported a study of the excited state dynamics of alkali metal ion–DB18C6 complexes in methanol,³ and on the basis of the UV absorption spectra they also suggested that K^+ forms the most rigid complex with DB18C6.

In the 1990s, metal ion–crown ether complexes were extensively investigated in the gas phase with mass spectrometric techniques^{4–16} and ion mobility methods^{17,18} and more recently by IR spectroscopy.^{19–25} The structure of benzo-crown ethers and their complexes with neutral guest species was studied by UV and IR spectroscopy under jet-cooled conditions.^{26–32} Very

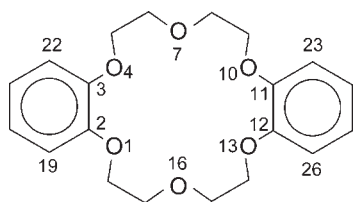
recently, Kim and co-workers reported UV spectra of metal ion–DB18C6 complexes produced by electrospray.^{33–35} Although they cooled the complexes to 150 K in the gas phase, the UV spectra still showed only broad features.

In the present work, we return to the original method of Pedersen for studying cation–DB18C6 complexes, UV spectroscopy, but we combine it with cooling in an ion trap.^{36,37} We report here UV photodissociation (UVPD) spectra of DB18C6 complexes with alkali metal ions, $\text{M}^+\cdot\text{DB18C6}$ ($\text{M} = \text{Li}, \text{Na}, \text{K}, \text{Rb}, \text{and Cs}$). The $\text{M}^+\cdot\text{DB18C6}$ complexes are produced by nanoelectrospray, mass-selected by a quadrupole mass spectrometer and irradiated with a UV laser in a cold, 22-pole ion trap, where the complexes are cooled to $\sim 10\text{ K}$.^{36,37} Isolating and cooling the complexes in the gas phase greatly simplify the UV spectra and provide well-resolved vibronic bands. Furthermore, we measure conformer-specific infrared spectra via IR-UV double resonance in the CH stretching ($2800\text{--}3120\text{ cm}^{-1}$) region. Although the observed IR spectra are more congested than predicted by quantum chemical calculations because of Fermi resonance interactions, we can still use these spectra to distinguish peaks in the UV spectra that belong to different conformers. The use of density functional theory (DFT) allows us to determine the conformation of the complexes.

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Scheme 1



2. EXPERIMENTAL AND COMPUTATIONAL METHODS

The details of our experimental approach have been given elsewhere.³⁸ Briefly, $M^+ \cdot \text{DB18C6}$ ($M = \text{Li}, \text{Na}, \text{K}, \text{Rb}, \text{and Cs}$) complexes are produced continuously at atmospheric pressure via nanoelectrospray using 20–200 μM solutions of DB18C6 and metal chloride salts in methanol. Ions are transferred into a vacuum chamber through an 11.5 cm long stainless steel capillary and injected into an ion funnel, which is used for collecting and focusing the divergent stream of ions exiting the capillary. After leaving the funnel, the ions pass through a hexapole ion guide and a quadrupole mass filter before being deflected 90° by an electrostatic quadrupole and stored in a room-temperature octopole ion trap. The trap accumulates the continuous stream of ions and ejects them in packets with a repetition rate of 10 Hz. The resulting ion packets are deflected by a second electrostatic quadrupole before entering a 22-pole RF ion trap, which is housed in a copper box and cooled by a closed-cycle He refrigerator to 4 K. The trapped ions are cooled internally and translationally to low temperature through collisions with cold He buffer gas, which is injected into the trap volume using a pulsed valve. The vibrational temperature of ions is estimated to be ~ 10 K.^{36,37} The trapped ions are then irradiated with a UV laser pulse, which causes some fraction of them to dissociate. The resulting charged photofragments, as well as the remaining parent ions, are released from the trap, deflected 90° by a third quadrupole deflector, mass-analyzed by a quadrupole mass filter, and detected with a channel electron multiplier. Ultraviolet photodissociation spectra of parent ions are obtained by plotting the yield of a particular photofragment ion as a function of the wavenumber of the UV laser. The UV light comes from the frequency-doubled output of a dye laser, which is pumped with the third harmonic of a Nd:YAG laser. For IR-UV double-resonance spectroscopy, we use a tunable IR optical parametric oscillator (OPO) pumped with a Nd:YAG laser. The output pulse of the IR OPO counter-propagates collinearly with the UV pulse, arriving in the trap ~ 100 ns prior to it. The wavenumber of the UV laser is fixed to a particular vibronic band in the UVPD spectra for monitoring the population of a conformer. The wavenumber of the OPO is scanned while monitoring fragment ion intensity induced by the UV laser. Conformer-specific IR spectra are obtained by plotting the yield of a particular photofragment as a function of the OPO wavenumber. For distinguishing vibronic bands due to different conformers, we measure UVPD spectra with the OPO wavenumber fixed to a particular vibrational band. The detailed procedure of IR-UV double-resonance spectroscopy is described in the Supporting Information.

For geometry optimization of the DB18C6 complexes with alkali metal ions, we first use a classical force field to find conformational minima. The initial conformational search is performed by using the mixed torsional search with low-mode sampling and the MMFF94s force field as implemented in MacroModel V. 9.1.³⁹ Within a 41.84 kJ/mol (10.00 kcal/mol) energy window, 219, 201, 123, 109, and 83 conformations are found for the Li^+ , Na^+ , K^+ , Rb^+ , and Cs^+ complexes, respectively. Minimum-energy conformers found with the force field calculations are then optimized at the M05-2X/6-31+G(d) level with loose optimization criteria with the GAUSSIAN09 program package.⁴⁰

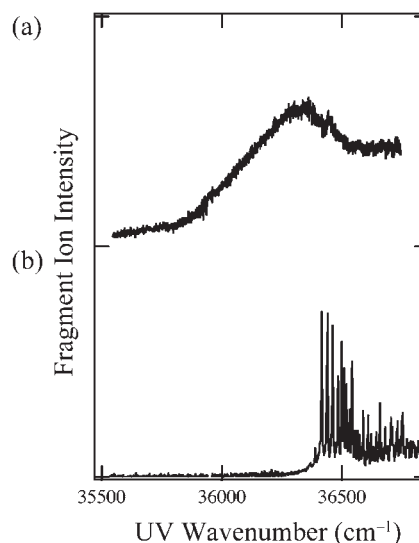


Figure 1. UVPD spectra of (a) uncooled $\text{K}^+ \cdot \text{DB18C6}$ and (b) $\text{K}^+ \cdot \text{DB18C6}$ that is cooled in the 22-pole ion trap. The temperature of the cooled complex is estimated to be ~ 10 K.

Unique minima obtained by comparison of relative energies (tolerance 0.05 kJ/mol) and rotational constants are further optimized at the M05-2X/6-31+G(d) and wB97XD/6-31++G(d,p) levels. Vibrational analysis is carried out for the optimized structures at the same level of calculation. The oscillator strength and transition energy of electronic transitions are obtained with time-dependent DFT (TD-DFT) at the M05-2X/6-31+G(d) level. In addition, geometry optimization in the S_1 state is also performed with TD-DFT at the M05-2X/6-31+G(d) level; the optimized structures in the S_0 state are used as the starting point for optimization of the excited state structures. For Rb and Cs, we use several effective core potentials (ECPs) such as LanL2DZ, Stuttgart RLC, Stuttgart RSC 1997, and CRENL; functions of these ECPs are obtained from a database of basis sets.⁴¹ For comparison with the calculated results of Choi et al.,³⁴ we also calculate all the complexes using the B3LYP functional. All the experimental and calculated results and their analysis are described in the Supporting Information. We use the character “S” for referring to figures in the Supporting Information, e.g., “Figure 1S”.

3. RESULTS

3.1. UV Spectroscopy. Figure 1 displays the UVPD spectra of (a) room-temperature $\text{K}^+ \cdot \text{DB18C6}$ and (b) $\text{K}^+ \cdot \text{DB18C6}$ that is cooled in the 22-pole ion trap. The uncooled complex has a broad absorption around 36 300 cm^{-1} , which is similar to the UVPD spectrum measured at 150 K.³⁴ The UV spectrum of the cooled $\text{K}^+ \cdot \text{DB18C6}$ complex consists of many sharp bands, with the band origin clearly observed at 36 415 cm^{-1} . From the intensity and the vibrational frequency of hot bands observed on the low-energy side of the band origin, we estimate the temperature of the cooled complex to be ~ 10 K. Comparison of the UV spectra in Figure 1 indicates that the broad absorption of the uncooled complex is mainly due to thermal congestion.

Figure 2 shows the UVPD spectra of the cooled $M^+ \cdot \text{DB18C6}$ ($M = \text{Li}, \text{Na}, \text{K}, \text{Rb}, \text{and Cs}$) complexes in the 35 500–38 000 cm^{-1} region together with the LIF spectrum of bare, jet-cooled DB18C6.²⁸ Except for the Li^+ complex, these UVPD spectra are measured by monitoring the yield of the M^+ photofragment ions. The Li^+ ion ($m/z = 7$) could not be detected with

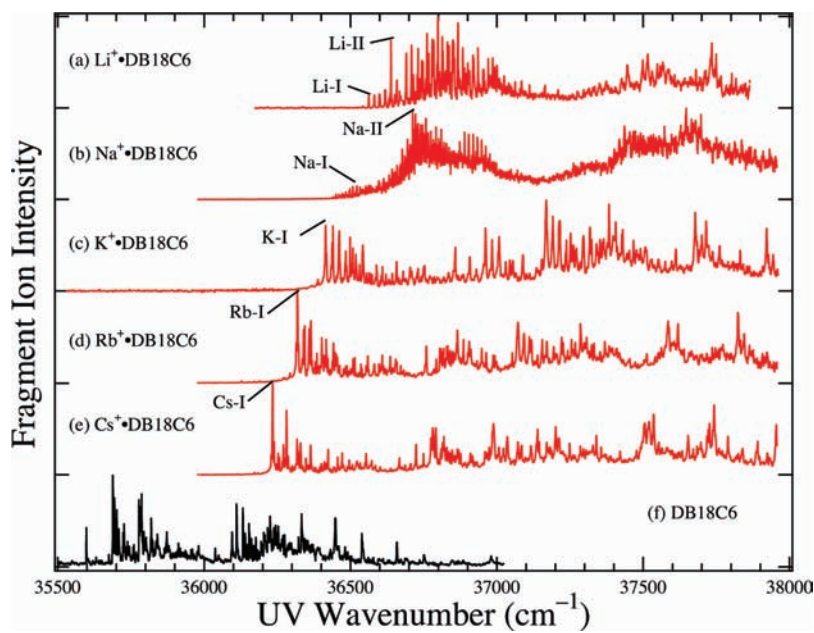


Figure 2. UVPD spectra of the $M^+ \cdot \text{DB18C6}$ ($M = \text{Li}, \text{Na}, \text{K}, \text{Rb}, \text{and Cs}$) complexes with the LIF spectrum of jet-cooled DB18C6 reported by Kusaka et al. (ref 28).

the second quadrupole mass filter because of its low-mass cutoff. However, upon UV excitation the $\text{Li}^+ \cdot \text{DB18C6}$ complex produces a variety of photofragments that result from the decomposition of the ether ring, as shown in Figure 1S. The spectrum in Figure 2a is measured by detecting the photofragment cation at $m/z = 309$, which corresponds to the loss of $\text{CH}_2\text{—OCH}_2\text{CH}_2$. All the $M^+ \cdot \text{DB18C6}$ complexes show a blue shift of the absorption relative to uncomplexed DB18C6. In UV spectroscopic studies of $\text{DB18C6}-(\text{H}_2\text{O})_n$ complexes, the attachment of H_2O molecules to the α oxygen atoms (O1, O4, O10, and O13, see Scheme 1) through the $\text{O—H} \cdots \text{O}$ hydrogen bond causes a blue-shift of the $S_1\text{—}S_0$ absorption.^{28,29} This suggests that the metal ions in the $M^+ \cdot \text{DB18C6}$ complexes are bonded to the α oxygen atoms, and the shift in the band origin in the UVPD spectra is an indication of the interaction strength. The Li^+ complex shows the most blue-shifted band origin, indicating the strongest interaction. This result is consistent with enthalpies of binding measured by collision-induced dissociation.¹⁶ For the Rb^+ and Cs^+ complexes, the band origin gradually shifts to the red with respect to the position of the K^+ complex. This indicates that the interaction between the α oxygen atoms and the metal ion becomes progressively weaker from K^+ to Cs^+ .

Figure 3 displays expanded views of the UVPD spectra of the $\text{K}^+ \cdot \text{DB18C6}$, $\text{Rb}^+ \cdot \text{DB18C6}$, and $\text{Cs}^+ \cdot \text{DB18C6}$ complexes around their respective band origins. In the previous study, the UV spectra of the K^+ , Rb^+ , and Cs^+ complexes exhibited only broad features.³⁴ In contrast, the cold UV spectra of Figure 3 exhibit many resolved features that provide a rich source of information on the complexes. As is highlighted by solid lines, the spectral features exhibit a band splitting of a few cm^{-1} . On the basis of the study of UV spectroscopy for jet-cooled DB18C6, we assign this structure to the exciton splitting that occurs between the benzene rings.²⁸ The doublet structure appears repeatedly in the spectra, built on top of vibrational progressions of ~ 20 and 47 cm^{-1} , which are due to low-frequency vibrations of the complexes (labeled α and β in Figure 3). In addition, the spectra

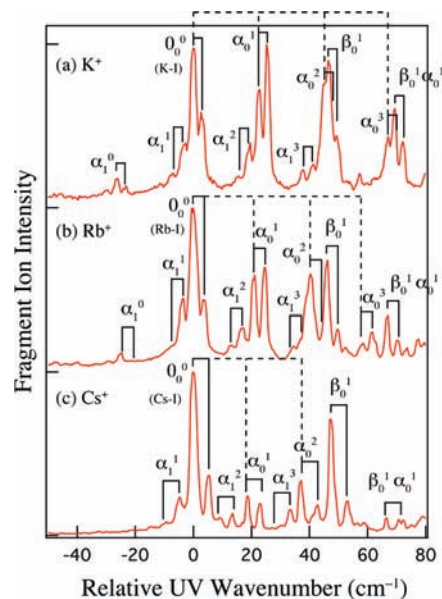


Figure 3. UVPD spectra of the DB18C6 complexes with K^+ , Rb^+ , and Cs^+ around their respective origin bands. The horizontal axis is the UV wavenumber relative to the band origin. The doublet structure connected by solid lines is attributed to the exciton splitting due to the intramolecular interaction between the benzene rings. The low-frequency vibrations appearing in these spectra are labeled by α and β . Superscript and subscript integers with α and β denote quantum numbers of these vibrations in the S_1 and S_0 states, respectively. Progressions of vibrations α from the origin band are presented with dotted lines.

have weak features assignable to hot bands. Superscript and subscript integers with α and β denote quantum numbers of these vibrations in the S_1 and S_0 states, respectively. The progression starts not only from the origin band (0_0^0) but also from a hot band (α_1^0). As seen in Figure 2 or Figure 8S, the

Table 1. Positions of the Band Origin (cm^{-1}), Exciton Splitting (cm^{-1}), and Vibrational Frequencies (cm^{-1}) Observed in the UVPD Spectra of the $\text{K}^+\cdot\text{DB18C6}$, $\text{Rb}^+\cdot\text{DB18C6}$, and $\text{Cs}^+\cdot\text{DB18C6}$ Complexes

	$\text{K}^+\cdot\text{DB18C6}$	$\text{Rb}^+\cdot\text{DB18C6}$	$\text{Cs}^+\cdot\text{DB18C6}$
band origin (cm^{-1})	36415	36319	36234
exciton splitting (cm^{-1})	2.7	4.0	5.3
mode α (S_0) (cm^{-1}) ^a	27	25	27
mode α (S_1) (cm^{-1}) ^a	23	21	19
mode β (cm^{-1}) ^a	47	46	47
mode 1 (cm^{-1})	754	754	754
mode 6 (cm^{-1})	546	547	547
mode 12 (cm^{-1})	968	967	967
mode A (cm^{-1}) ^a	903	903	905
mode B (cm^{-1}) ^a	1263	1267	1271
band origin (cm^{-1}) ^b	36375 ± 7	36268 ± 10	36173 ± 12

^a Unidentified normal modes are labeled with α , β , A, and B. Normal modes A and B are due to vibrations of the benzene rings. ^b Ref 34.

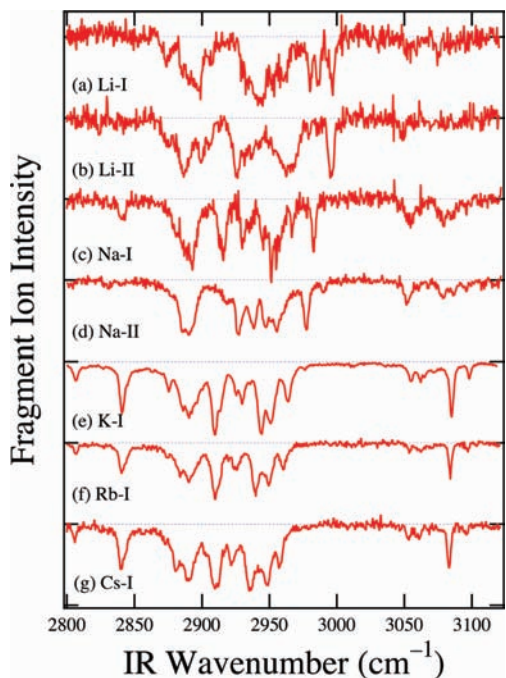


Figure 4. IR-UV double-resonance spectra of the $\text{M}^+\cdot\text{DB18C6}$ ($\text{M} = \text{Li}, \text{Na}, \text{K}, \text{Rb}, \text{and Cs}$) complexes in the CH stretching region.

vibronic pattern around the origin band appears also in the higher-frequency region. These patterns are due to vibrations involving mainly the benzene rings, and the frequency is almost the same for the K^+ , Rb^+ , and Cs^+ complexes. The intervals of the exciton splitting and vibrational frequencies are collected in Table 1.

3.2. IR Spectroscopy. All the UVPD spectra in Figure 2 have a number of vibronically resolved bands. To distinguish vibronic bands due to different conformers, IR-UV double-resonance spectra are measured using strong vibronic bands. Figure 4 shows the IR-UV spectra of the $\text{M}^+\cdot\text{DB18C6}$ complexes in the CH stretching ($2800\text{--}3120\text{ cm}^{-1}$) region. For the Li^+ and Na^+ complexes, two kinds of IR spectra are found; the IR-UV

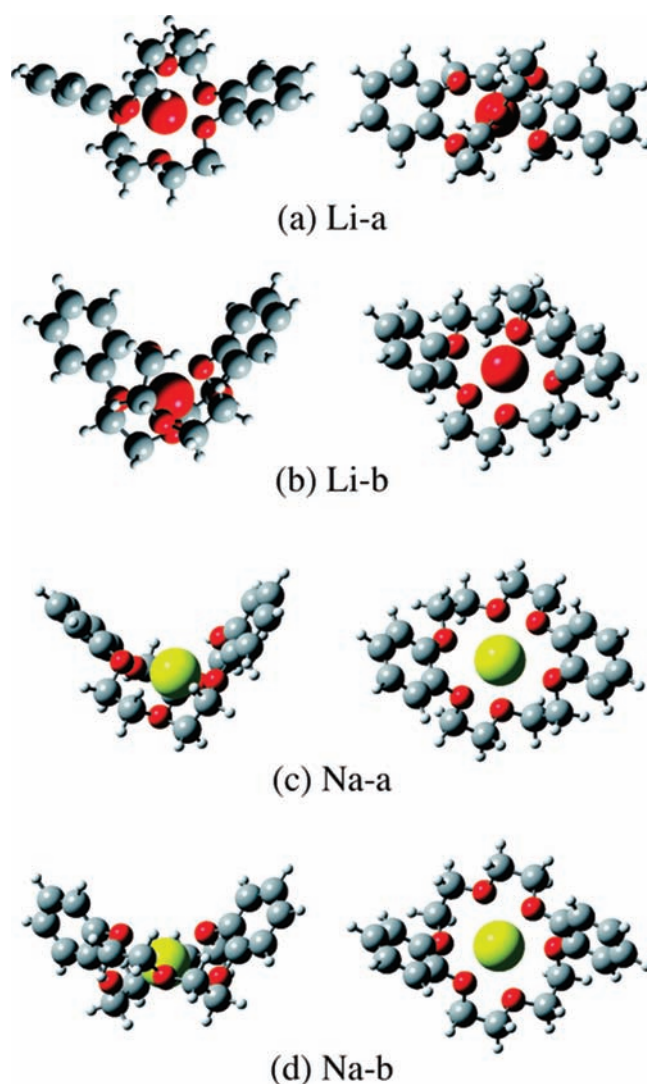


Figure 5. Structure of (a, b) $\text{Li}^+\cdot\text{DB18C6}$ and (c, d) $\text{Na}^+\cdot\text{DB18C6}$ complexes optimized at the M05-2X/6-31+G(d) level of theory.

spectra in Figures 4a–d are measured by monitoring the intensity of the vibronic bands labeled Li-I, Li-II, Na-I, and Na-II in Figure 2. These results indicate that there exist at least two conformers for the Li^+ and Na^+ complexes. The IR-UV spectra of the K^+ , Rb^+ , and Cs^+ complexes in Figures 4e–g are observed by monitoring the origin band (K-I, Rb-I, and Cs-I in Figure 2). Using other vibronic bands of these complexes produces the same IR spectra, demonstrating that there is only one stable isomer for each of them. The similarity of the IR spectra of the K^+ , Rb^+ , and Cs^+ complexes implies that the conformation of DB18C6 in each of them will be similar. We also perform hole-burning experiments with the IR laser to confirm the number of conformers. UVPD spectra of the $\text{M}^+\cdot\text{DB18C6}$ complexes are measured under the IR irradiation with the wavenumber fixed to a particular vibrational band. With these UV measurements, we distinguish which peaks in the UVPD spectra belong to different conformers (see Figures 3S and 4S). If different conformers have a similar IR absorption spectrum, it is difficult to distinguish them with IR spectroscopy. Therefore, the number of conformers determined in this study is a lower limit. Nevertheless, we have

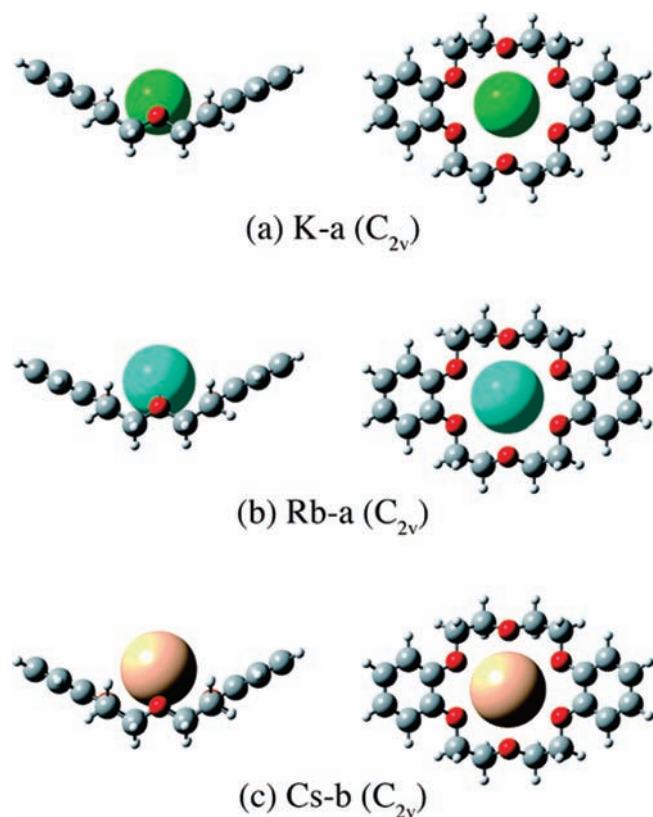


Figure 6. Structure of (a) $K^+ \cdot DB18C6$, (b) $Rb^+ \cdot DB18C6$, and (c) $Cs^+ \cdot DB18C6$ complexes optimized at the M05-2X/6-31+G(d) level of theory.

carefully measured IR spectra for most of the strong vibronic bands in the UVPD spectra, and all the IR spectra can be classified into one of the spectra in Figure 4. On this basis, we can say that if any other conformers exist under the cold conditions of our ion trap they represent at best a minor fraction of the total complex population.

After sorting the vibronic bands by conformers, we can analyze the vibronic structure in the UVPD spectra more definitively, as described in detail in the Supporting Information. For the Li^+ complex, the origin band of the two stable conformers is strong, one of which is accompanied by an extensive low-frequency progression with an interval of 19.1 cm^{-1} (see Figure 5S). For the Na^+ complex, the origin band is too weak to be identified, and low-frequency progressions are extensive for both of the stable conformers (see Figure 6S). As seen in Figure 3, low-frequency progressions can be seen also for the K^+ , Rb^+ , and Cs^+ complexes, but these are not so extensive. We discuss more fully below the relation between the activity of the low-frequency progressions in the UV spectra and the conformation of the complexes.

4. DISCUSSION

4.1. Determination of the Conformation for the $M^+ \cdot DB18C6$ Complexes. We determine the structure of the $M^+ \cdot DB18C6$ complexes with the aid of DFT and TD-DFT. For the Li^+ and Na^+ complexes, the two conformers identified in the UVPD spectra can be attributed to the most and the second most stable forms calculated at the M05-2X level because these are substantially more stable than others. (The calculated total

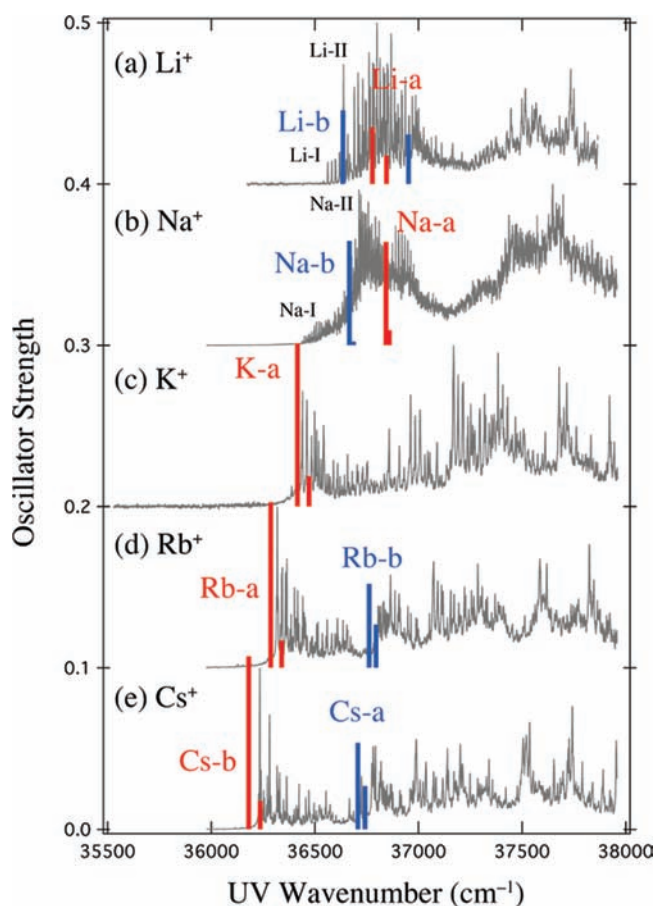


Figure 7. Oscillator strengths of electronic transitions for $M^+ \cdot DB18C6$ ($M = Li, Na, K, Rb, Cs$) calculated with TD-DFT at the M05-2X/6-31+G(d) level and the Stuttgart RLC potential for Rb and Cs. Calculated transition energies are scaled with a scaling factor of 0.8340 to compare the calculated and experimental spectra. The scaling factor is determined to reproduce the transition energy of the origin band of $K^+ \cdot DB18C6$.

energy for all the conformers is described in the Supporting Information.) Figure 5 displays the lowest-energy structures of the $Li^+ \cdot DB18C6$ and $Na^+ \cdot DB18C6$ complexes calculated at the M05-2X/6-31+G(d) level. In the previous study by Choi et al., the $Li^+ \cdot DB18C6$ complex was ascribed to another structure (Li-c in Figure 9S) on the basis of B3LYP calculations.³⁴ However, the possibility of the structure (Li-c) as the conformation of $Li^+ \cdot DB18C6$ is ruled out by the absence of the exciton splitting in the UVPD spectrum, as described in the Supporting Information.

Calculated structures of the K^+ , Rb^+ , and Cs^+ complexes are shown in Figure 6. The K^+ complex has a boat-type C_{2v} conformer (K-a in Figure 6a). This is the most stable conformer at all the levels of B3LYP/6-31+G(d), M05-2X/6-31+G(d), and wB97XD/6-31++G(d,p). From the similarities among the UVPD spectra and the IR-UV spectra, we expect the Rb^+ and Cs^+ complexes to have a C_{2v} structure similar to K-a. However, the predominance of the boat C_{2v} conformers (Rb-a and Cs-b in Figure 6) is not so clear from the calculations of the Rb^+ and Cs^+ complexes. For the Rb^+ complex, all the calculations show that the boat conformation (Rb-a) is the most stable form, but there is another conformer (Rb-b in Figure 12S) with a total energy very close to that of Rb-a. In the case of the Cs^+ complex, the boat C_{2v}

Table 2. Distance between the Center of the Benzene Rings and the Difference in the Distance between the S_0 and S_1 States Determined with Geometry Optimization at the M05-2X/6-31+G(d) Level and Stuttgart RLC Potential

conformers	S_0 (Å)	S_1 (Å)	difference (Å)
Li-a	7.87	7.94	0.07
Li-b	6.39	6.80	0.41
Na-a	7.48	7.98	0.50
Na-b	7.72	8.18	0.46
K-a	9.09	9.11	0.01
Rb-a	9.24	9.36	0.12
Cs-b	9.25	9.37	0.12

conformer (Cs-b) is the second most stable form at the M05-2X and wB97XD calculations; the most stable one (Cs-a in Figure 13S) has a structure similar to Rb-b. The stabilization energy is dependent on the DFT functionals and the ECPs used in geometry optimization. Therefore, the structure of the Rb^+ and Cs^+ complexes cannot be determined definitely on the basis of their calculated energy. As shown in Figure 23S, the observed IR spectra are more congested than predicted by the DFT calculations, which is due to the Fermi resonance interactions between the CH stretching and overtone of the CH bending vibrations. Therefore, comparison of the observed and calculated IR spectra does not provide a definitive conclusion for the complex structure.

In contrast, the assignment of the structure is confirmed with TD-DFT calculations, which do not strongly depend on the calculation methods. Figure 7 shows electronic spectra calculated with TD-DFT at the M05-2X/6-31+G(d) level. For the calculated electronic transition energies, a scaling factor of 0.8340 is used for comparison with the experimental spectra. This factor is determined to reproduce the transition energy of the origin band of the $K^+ \cdot DB18C6$ complex with the most stable conformer (K-a). The transition energy calculated for Rb-a (C_{2v}) and Cs-b (C_{2v}) is coincident with the position of the origin bands for the Rb^+ and Cs^+ complexes, while Rb-b and Cs-a show band origins that are blue-shifted with respect to that of K-a rather than red-shifted (as observed). These results reinforce the assignment of the structure of the Rb^+ and Cs^+ complexes to Rb-a and Cs-b. In the case of the Li^+ and Na^+ complexes, the second most stable conformers (Li-b and Na-b) show electronic transition energies lower than those of the most stable ones (Li-a and Na-a). Therefore, we attribute the vibronic bands of Li-I and Li-II to conformers Li-b and Li-a, respectively. Similarly, Na-I and Na-II in the UVPD spectrum are ascribed to Na-b and Na-a, respectively. In both Li-a and Li-b, DB18C6 distorts its ether ring to hold the Li^+ ion because the cavity size of DB18C6 is much larger than Li^+ . This distortion is also seen for Na-a and Na-b. However, the distance between the oxygen atoms and the Na^+ ion is longer (2.39 Å) than that of Li-a and Li-b (2.09 Å), and the ether ring is not so highly distorted as in the Li^+ complex. In K-a, Rb-a, and Cs-b, the ether ring of DB18C6 opens the most, adopting a boat (C_{2v}) conformation. The K^+ ion is in the center of the cavity, while the Rb^+ and Cs^+ ions sit on top of the open cavity because their ion diameters are larger than the cavity size. The conformer structures in Figures 5 and 6 are similar to those predicted by Bowers and co-workers on the basis of ion mobility measurements for alkali complexes with B18C6.^{17,18} Very recently, Rodriguez and Lisy investigated the structure of $M^+(18C6)(H_2O)_{1,2}$

($M = Li, Na, K, Rb,$ and Cs) by IR predissociation spectroscopy of $M^+(18C6)(H_2O)_{1,2}Ar$ and DFT calculations.⁴² Under hydrated conditions, the 18C6 ring adopts not distorted but open forms, encapsulating the Li^+ and Na^+ ions inside the cavity with an H_2O molecule. For the K^+ ion, the “best fit” with the 18C6 cavity weakens the $K^+ \cdots H_2O$ interaction. In the case of the Rb^+ and Cs^+ complexes, the metal ions are positioned above the plane of the 18C6 ring. The sharp structural difference between Li, Na, and K and the general trends for K through Cs of the $M^+(18C6)(H_2O)_{1,2}$ complexes are similar to those of the $M^+(DB18C6)$ clusters.

We can compare the gas-phase structures in Figures 5 and 6 with those obtained by X-ray diffraction analysis of alkali metal salts with DB18C6. The Cambridge Structural Database (CSD ver. 5.31) of the Cambridge Crystallographic Data Centre⁴³ contains 32, 59, 10, and 3 crystal structures for the Na^+ , K^+ , Rb^+ , and Cs^+ complexes, respectively. No crystal data of the $Li^+ \cdot DB18C6$ complex were put into the database. Rb^+ and Cs^+ ions are each held by two DB18C6 molecules in all the crystals, implying that they are too large to be located inside the DB18C6 cavity. For all the K^+ complexes in crystals, the DB18C6 part has a boat-type open form, similar to that in the gas phase (Figure 6a).^{44,45} In contrast, the $Na^+ \cdot DB18C6$ complex in crystals has structures quite different from that in the gas phase.^{46–49} In crystals, the conformation of the DB18C6 adopts an open form, similar to that in the $K^+ \cdot DB18C6$ complex, whereas the ether ring is distorted in the gas phase to fitting the cavity size to Na^+ (Figures 5c and 5d). In crystals, counterions also take part in the coordination of metal cations, which significantly influences the complex structure. Since alkali metal complexes of DB18C6 are separated from counteranions in solution, the structure of the complexes in solution should be more similar to that in the gas phase than in crystals.

4.2. Conformation Change upon UV Irradiation. Activity of low-frequency vibrational progressions in a UV spectrum is an indication of a structural change along the vibrational coordinates of low-frequency normal modes upon S_1-S_0 excitation. Geometry optimization in the S_1 and S_0 states well explains the activity of the progressions observed in the UVPD spectra of the complexes we have measured. Vibrational analysis in the S_0 state shows that all the conformers in Figures 5 and 6 have low-frequency vibrations of $\sim 20 \text{ cm}^{-1}$, which are torsional motions of the ether ring. On the basis of these results, the low-frequency progressions observed in the UVPD spectra can be attributed to torsional motions of the ring in the S_1 state. As shown in Figure 7S, the torsional motions involve a change of the distance between the benzene rings. Therefore, the distance can be used as a measure of the structural change upon S_1-S_0 excitation. Table 2 shows the distance between the benzene rings in the S_0 and S_1 states and the difference in the distance for the conformers in Figures 5 and 6. Li-b, Na-a, and Na-b show much larger differences ($\sim 0.45 \text{ Å}$) than those of Li-a, K-a, Rb-a, and Cs-b ($\sim 0.10 \text{ Å}$). These calculated results coincide well with the experimental observations in that one of the Li^+ conformers and both of the Na^+ conformers show the extensive and intense low-frequency vibrational progressions, as seen in Figures 5S and 6S. All the conformers in Table 2 show an increase of the distance between the benzene rings from the S_0 to the S_1 state. The conformers in which the distance in the S_0 state is shorter than $\sim 7.80 \text{ Å}$ tend to have larger structural change, which causes the extensive low-frequency progressions in the UVPD spectra. This indicates that complexes that are highly distorted open the ether

ring more extensively upon S_1 – S_0 excitation, especially in the cases of Li^+ and Na^+ . As a result, the metal ions will be exposed more on the outside of the ether ring, which may change their chemical activity. This suggests that the encapsulation of metal ions in the cavity of DB18C6 could perhaps be controlled by UV irradiation.

5. SUMMARY

We have measured UV photodissociation (UVPD) and IR–UV double-resonance spectra of the $M^+ \cdot DB18C6$ ($M = Li, Na, K, Rb,$ and Cs) complexes produced by nanoelectrospray and cooled to ~ 10 K in the 22-pole ion trap. All the $M^+ \cdot DB18C6$ complexes show a number of well-resolved vibronic bands. With the aid of quantum chemical calculations, we conclude the following: (1) The $Li^+ \cdot DB18C6$ and $Na^+ \cdot DB18C6$ complexes exhibit two stable conformers in our cold ion trap, whereas the K^+, Rb^+ , and Cs^+ complexes show only one. (2) For holding Li^+ or Na^+ , DB18C6 distorts the ether ring to fit the cavity size to the ion diameter. (3) In the $K^+ \cdot DB18C6$ complex, the DB18C6 opens the ether cavity the most, and the K^+ ion fits well inside. (4) For the Rb^+ and Cs^+ complexes, DB18C6 has an open form similar to the $K^+ \cdot DB18C6$ complex. However, since the ion size is too large to fit inside the cavity of DB18C6, the ion sits on top of the ether ring, and the interaction between the α oxygen atoms and the ions becomes weaker than that of the K^+ complex. (5) Complexes that are highly distorted to hold metal ions open the ether ring largely upon S_1 – S_0 excitation. As a result, the encapsulation of metal ions in the cavity of DB18C6 could perhaps be controlled by the UV irradiation.

■ ASSOCIATED CONTENT

S Supporting Information. Results of IR–UV double-resonance spectroscopy, analysis of the vibronic structure in the UVPD spectra, optimized structures, electronic spectra calculated with the TD–DFT calculations, and a full list of authors of ref 40. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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